Efficient computational modeling of piano strings for real-time synthesis using mass-spring chains, coupled finite differences, and digital waveguide sections

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> > ASA-2010

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Goals

• Goals

• Outline

String Synth History

Mass-Spring Model

Finite Differences

- Overall Goal: Ultimate Virtual Piano
- Current Focus: Audibly Perfect Piano-String Synthesis
- Method: Start with High Accuracy, then Simplify





Outline

Goals

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String Synth History

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Finite Differences

- Brief History of Virtual Strings
 - Ideal Strings
 - Stiff Piano Strings
 - Nonlinear Piano Strings
- General Mass-Spring String Model
- Finite Difference Implementation





String Synth History Mass-Spring Model Finite Differences

Brief History of String Sound Synthesis





D'Alembert's PDE for the Ideal String (1747)

- Goals
- Outline

String Synth History

- Ideal String
- Traveling Waves
- Digital Waveguide
- Plucked String
- Struck String
- Stiff Strings
- Stiff Waveguides
- Commuted Stiffness
- Sound Example
- Nonlinear Coupling
- Coupling Effects
- Nonlinear Synthesis
- Synthesis
- Zooming Out

Mass-Spring Model

Finite Differences



Wave equation PDE is derived from Newton's second law (f = ma) by equating Brook Taylor's "restoring force" Ky'' (1713) to mass-density times acceleration $\epsilon \ddot{y}$





D'Alembert's Traveling-Wave Solution (1747)

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Finite Differences

 In the same paper, d'Alembert also showed that the ideal-string PDE was satisfied by *traveling waves* in either direction:

$$y(t,x) = y_r\left(t - \frac{x}{c}\right) + y_l\left(t + \frac{x}{c}\right)$$

where

$$c = \sqrt{\frac{K}{\epsilon}} =$$
 wave propagation speed

 These ideas were developed into essentially modern form by Euler (1707–1783)





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Finite Differences

Digital Waveguide Models (1985)

We sample d'Alembert's traveling-wave components to make *digital waveguide models.*

Lossless digital waveguide $\stackrel{\Delta}{=}$

bidirectional delay line

at some wave impedance $R = \sqrt{K\epsilon}$



Useful for *efficient* models of strings, bores, plane waves, conical waves, and more





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Finite Differences

Ideal Plucked String (Displacement Waves)

Digital waveguide string models are *extremely efficient* computationally [$\mathcal{O}(1)$ complexity per sample of output]



- Load each delay line with *half* of initial string displacement
- Sum of upper and lower delay lines = string displacement
- Insert a linear, time-invariant filter in the loop for any desired *attenuation* and *dispersion* as a function of frequency





Ideal Struck String (Velocity Waves)



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Finite Differences



Velocity waves v easily converted to *force* waves f which are proportional to *string slope* y':

$$-Ky' \stackrel{\Delta}{=} f = f^+ + f^- = Rv^+ - Rv^-$$



ASA-2010 - 9 / 31



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- Commuted Stiffness
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Mass-Spring Model

Finite Differences

Plucked/Struck Sound Example (1988)

Bach A-Minor Concerto—Orchestra Part: (WAV) (MP3)

 Executed in real time on one Motorola DSP56001 (20 MHz clock, 128K SRAM)





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- Outline

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- Stiff Strings
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- Commuted Stiffness
- Sound Example
- Nonlinear Coupling
- Coupling Effects
- Nonlinear Synthesis
- Synthesis
- Zooming Out

Mass-Spring Model

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Plucked/Struck Sound Example (1988)

 Developed for the NeXT Computer introduction at Davies Symphony Hall, San Francisco, 1988





- Goals
- Outline

- Ideal String
- Traveling Waves
- Digital Waveguide
- Plucked String
- Struck String
- Stiff Strings
- Stiff Waveguides
- Commuted Stiffness
- Sound Example
- Nonlinear Coupling
- Coupling Effects
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- Synthesis
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Plucked/Struck Sound Example (1988)

Bach A-Minor Concerto—Orchestra Part: (WAV) (MP3)

 Executed in real time on one Motorola DSP56001 (20 MHz clock, 128K SRAM)

• Developed for the NeXT Computer introduction at Davies Symphony Hall, San Francisco, 1988

 Solo violin part was played live by Dan Kobialka of the San Francisco Symphony





Stiff Strings

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Finite Differences

• Piano strings are *stiff*, resulting in significant *dispersion*:

$$\epsilon \ddot{y} = K y'' - \kappa y'''$$
 (Stiff string PDE)

The new "stiffness term" is proportional to

$$y'''' \stackrel{\Delta}{=} \frac{\partial^4}{\partial x^4} y(t, x) \qquad \Rightarrow$$

- Faster wave propagation at higher frequencies (hence the "dispersion" of traveling-wave shapes)
 - Provided in digital waveguide strings using an *allpass filter* having a delay that decreases with frequency (typically on the order of 10 poles and zeros)
- A spring-resistance to corner formation in the string (usually left out in string synthesis models!)
 See, e.g., Cremer 1984





Digital Waveguide Model for Stiff Strings

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Example waveguide string, three spatial-samples long:



 $H_a(z) =$ allpass giving desired delay vs. frequency for one sample





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Finite Differences

In practice, we pull out one or more samples of pure delay:

Consolidated Allpasses



- $zH_a^3(z)$ (or $z^2H_a^6(z)$) is designed as a single allpass filter
- ≈ 10 poles and zeros can handle hundreds of samples of dispersion with perceptual equivalence
- Not valid for distributed nonlinear behavior (discussed later)



ASA-2010 - 13 / 31



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Finite Differences

• Piano: (WAV) (MP3) (Stanford Sondius Project, ca. 1995)

Digital Waveguide Piano Sound Example





Digital Waveguide Piano Sound Example

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Finite Differences

- Piano: (WAV) (MP3) (Stanford Sondius Project, ca. 1995)
- Uses the commuted synthesis technique





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- Outline

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Finite Differences

• Vertical excitation by hammer *couples nonlinearly* to the *longitudinal* direction:¹

$$\epsilon\ddot{\xi} = ES\xi'' + \frac{1}{2}ES[(y')^2]'$$

(longitudinal wave PDE)

where $\xi = \textit{longitudinal displacement}$

• Nonlinear coupling driven by

Nonlinear Piano Strings

$$[(y')^2]' \stackrel{\Delta}{=} \frac{\partial}{\partial x} \left[\frac{\partial}{\partial x} y(t, x) \right]^2$$

- Slope of square of string slope drives longitudinal waves
- Coupling from longitudinal back to transverse is smaller and typically *neglected*
- ¹E.g., Morse & Ingard 1968





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Finite Differences

Effects of Nonlinear Mode Coupling in Piano Strings

Main *audible effects* of nonlinear transverse-to-longitudinal coupling:

- Initial *longitudinal attack pulse* (the initial "shock noise" audible in a piano tone)
- 2. Inharmonic longitudinal modes
- 3. "Phantom partials"

(ongoing intermodulation products from transverse partials)

See Conklin lecture in "Five Lectures on the Acoustics of the Piano," edited by Anders Askenfelt (and listen to the sound examples)





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Finite Differences

Nonlinear Piano-String Synthesis

If one-way coupling is accurate enough (transverse to longitudinal but *not* vice versa), we can model its effects *separately* based on *observations* of transverse waves (Bank and Sujbert, JASA 2005):

- Longitudinal modes implemented as second-order resonators ("modal synthesis")
- Slope of squared slope *projected* onto longitudinal modes

Up-to-date summary:

"A Modal-Based Real-Time Piano Synthesizer" Balázs Bank, Stefano Zambon, and Federico Fontana IEEE Trans. Audio, Speech, and Language Processing May, 2010 (IEEE-ASLP)





Synthesis Strategies

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- Outline

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Mass-Spring Model

Finite Differences

Model complexity grows with dynamic level:

- 1. Linear Superposition: Transverse and longitudinal waveguides *decouple* into separate modes
- 2. Transverse \rightarrow Longitudinal Coupling
- 3. Transverse \leftrightarrow Longitudinal Coupling





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Finite Differences

Model complexity grows with dynamic level:

- 1. Linear Superposition: Transverse and longitudinal waveguides *decouple* into separate modes
- 2. Transverse \rightarrow Longitudinal Coupling
- 3. Transverse \leftrightarrow Longitudinal Coupling
- Proposed Synthesis Strategy:
 - Initial striking force determines the starting regime (1, 2, or 3)
 - Maximum slope |y'| over one or more periods indicates regime transitions
 - String model simplifies as it decays





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Finite Differences

The preceding involved several approximations:

- Neglected terms in PDEs
- Simplified synthesis models





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Finite Differences

The preceding involved several approximations:

- Neglected terms in PDEs
- Simplified synthesis models
- The following questions naturally arise:
 - How do we know for sure our approximations are inaudible?
 - We can listen, but could we miss an audible effect?
 - Could a difference become audible after more listening?





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What we really want is a *truth reference*—an "exact model"!





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The preceding involved several approximations:

- Neglected terms in PDEs
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- The following questions naturally arise:
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- There are software tools (*e.g.*, from *perceptual audio coding*) for measuring the *audible equivalence* of two sounds:
 - "Original" and "Encoded" for a CODEC
 - "Exact" and "Computationally Efficient" for piano models





String Synth History Mass-Spring Model Finite Differences







- Goals
- Outline

Mass-Spring Model

- 3D Mass-Springs
- Mass-Spring String
- Adding Stiffness
- Shear Stiffness
- Three-Spring Design
- Shear Springs

Finite Differences





• Hooke's Law:

 $||\underline{f}_1|| = k \cdot |l_1 - l_0| \qquad (l_0 = \text{spring rest length})$

• Vector Equation of Motion $(\underline{f}_i \in \mathbf{R}^3, \underline{x}_i \in \mathbf{R}^3)$:

$$\underline{f}_1 = k \cdot (\|\underline{x}_2 - \underline{x}_1\| - l_0) \cdot \frac{\underline{x}_2 - \underline{x}_1}{\|\underline{x}_2 - \underline{x}_1\|}$$

$$= k \left[1 - \frac{l_0}{\|\underline{x}_2 - \underline{x}_1\|} \right] (\underline{x}_2 - \underline{x}_1) = m_1 \underline{\ddot{x}}_1$$

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ASA-2010 - 21 / 31

Mass-Spring String Model



where

$$\alpha_i \stackrel{\Delta}{=} k \cdot \left[1 - \frac{l_0}{\|\underline{x}_{i+1} - \underline{x}_i\|} \right]$$





Adding Stiffness

- Goals
- Outline

String Synth History

Mass-Spring Model

- 3D Mass-Springs
- Mass-Spring String
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- Shear Stiffness
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Finite Differences





Need Shear Stiffness Too



- Goals
- Outline

String Synth History

Mass-Spring Model

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Finite Differences







- Goals
- Outline

Mass-Spring Model

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Finite Differences



Three-Spring Design (Before Shear Springs Added)







Shear Springs

- Goals
- Outline

String Synth History

Mass-Spring Model

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Finite Differences







String Synth History Mass-Spring Model Finite Differences

Finite Difference Implementation





Digitizing the Flexible String

• Goals

• Outline

String Synth History

Mass-Spring Model

Finite Differences

- Digitization
- State-Space Form
- Summary

Recall:



$$\underbrace{f}_{i} = \alpha_{i} \cdot \left(\underline{x}_{i+1} - \underline{x}_{i} \right) + \alpha_{i-1} \cdot \left(\underline{x}_{i-1} - \underline{x}_{i} \right)$$

$$= \alpha_{i-1} \underline{x}_{i-1} - \left(\alpha_{i-1} + \alpha_{i} \right) \underline{x}_{i} + \alpha_{i} \underline{x}_{i+1}$$

where

$$\alpha_i \stackrel{\Delta}{=} k \cdot \left[1 - \frac{l_0}{\|\underline{x}_{i+1} - \underline{x}_i\|} \right]$$



Smith, Kuroda, Perng, Van Heusen, Abel



Equations of Motion, State-Space Form

- Goals
- Outline

String Synth History

Mass-Spring Model

- **Finite Differences**
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$$\begin{split} \underline{f}_{1} &= m_{1} \underline{\ddot{x}}_{1} = \alpha_{1} \cdot (\underline{x}_{2} - \underline{x}_{1}) \\ \vdots &\vdots \\ \underline{f}_{i} &= m_{i} \underline{\ddot{x}}_{i} = \alpha_{i-1} \underline{x}_{i-1} - (\alpha_{i-1} + \alpha_{i}) \underline{x}_{i} + \alpha_{i} \underline{x}_{i+1} \\ \vdots &\vdots \\ \underline{f}_{M} &= m_{M} \underline{\ddot{x}}_{M} = \alpha_{M-1} \cdot (\underline{x}_{M-1} - \underline{x}_{M}) \\ \\ \text{or } (3M \times 1): \end{split}$$

 $\underline{F} = \mathbf{M} \underline{\ddot{X}} = \mathbf{A} \underline{X}$

Digitization (Explicit Finite Difference Scheme):

$$\underline{X}_{n+1} = \left[2\mathbf{I} + \mathbf{M}^{-1}\mathbf{A} \right] \underline{X}_n - \underline{X}_{n-1} + B\underline{u}_n$$

where \underline{u}_n is the external input vector, and digitization is based on $\underline{\ddot{x}}_n \stackrel{\Delta}{=} \underline{x}_{n+1} - 2\underline{x}_n + \underline{x}_{n-1}$





Results to Date

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Finite Differences

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Presented this morning by Junji Kuroda





Summary

Goals

Outline

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Finite Differences

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• String-Synthesis History Reviewed

- Ideal Strings
- Stiff Piano Strings
- Nonlinear Piano Strings
- Mass-Spring Model = Accurate Benchmark Reference
 - $\circ\,$ Accuracy \sim time/space sampling density
 - Approximations can be quantified in practical cases
 - Given faster and more parallel computing, maybe we'll just go ahead and use it for real-time sound synthesis from performance!

