

# MUS420 Lecture

## Digital Waveguide Modeling of Bowed Strings

Julius O. Smith III ([jos@ccrma.stanford.edu](mailto:jos@ccrma.stanford.edu))  
Center for Computer Research in Music and Acoustics (CCRMA)  
Department of Music, Stanford University  
Stanford, California 94305

February 5, 2019

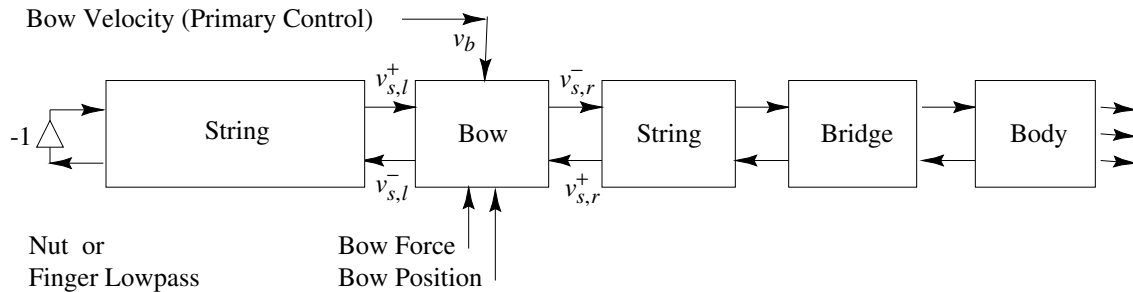
### Outline

- Physical Model
- Digital Waveguide Model
- Bow-String Theory and Computational Model
- Commuted Synthesis of Bowed Strings

# Bowed Strings

---

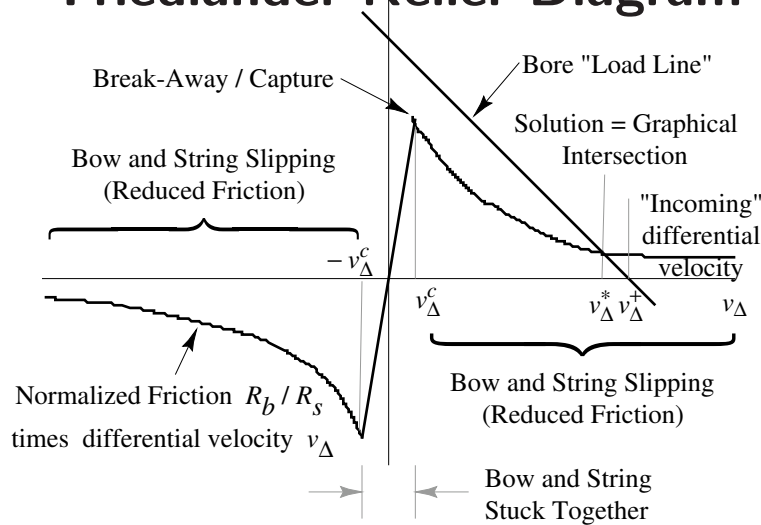
## Schematic Model



A schematic model for bowed-string instruments.

- Bow divides string into two sections
- Bow junction = nonlinear *two-port*
- Primary control variable = bow velocity  
⇒ *velocity waves* = natural choice of wave variable
- Bow-string interface is analogous to the reed-bore interface: Find intersection of bow-string friction curve with the string wave impedance “load line”
- In other words, a velocity input (injected equally to left and right) must be found such that the transverse force of the bow against the string is balanced by the reaction force of the moving string

## Friedlander-Keller Diagram



Overlay of normalized bow-string friction curve  $R_b(v_\Delta)/R_s$  with the bore “load line”  $v_\Delta^+ - v_\Delta$ . The “capture” and “break-away” differential incoming velocity is denoted  $v_\Delta^c$ . Note that increasing the bow force increases  $v_\Delta^c$  as well as enlarging the maximum force applied (at the peaks of the curve).

Applied Force = Friction Curve  $\times$  Differential Velocity  
 Reaction Force = String Wave Impedance  $\times$  Velocity Change

- Nominally,  $R_b(v_\Delta)$  is constant for  $|v_\Delta| \leq v_\Delta^c$ , where  $v_\Delta^c$  is both the capture and break-away differential velocity. (*static* coeff. of friction)
- For  $|v_\Delta| > v_\Delta^c$ ,  $R_b(v_\Delta)$  falls quickly to a low *dynamic* coefficient of friction
- Dynamic coefficient decreases with differential velocity

## Bow-String Scattering Junction

Friedlander-Keller diagram is solved when

$$R_b(v_\Delta) \times v_\Delta = R_s [v_\Delta^+ - v_\Delta]$$

which implies (in a manner analogous to the single reed case)

$$\begin{aligned} v_{s,r}^- &= v_{s,l}^+ + \hat{\rho}(v_\Delta^+) \cdot v_\Delta^+ \\ v_{s,l}^- &= v_{s,r}^+ + \hat{\rho}(v_\Delta^+) \cdot v_\Delta^+, \quad \text{where} \end{aligned}$$

$v_{s,r}$  = transverse string velocity on the *right* of the bow

$v_{s,l}$  = string velocity *left* of the bow ( $v_{s,l} = v_{s,r}$ )

$v_\Delta^+ \triangleq v_b - (v_{s,r}^+ + v_{s,l}^+) =$  “incoming differential velocity”

$v_b$  = bow velocity, and

$$\hat{\rho}(v_\Delta^+) = \frac{r(v_\Delta(v_\Delta^+))}{1 + r(v_\Delta(v_\Delta^+))} \quad \text{where}$$

$$r(v_\Delta) = 0.25R_b(v_\Delta)/R_s$$

$v_\Delta = v_b - v_s$       bow velocity minus string velocity

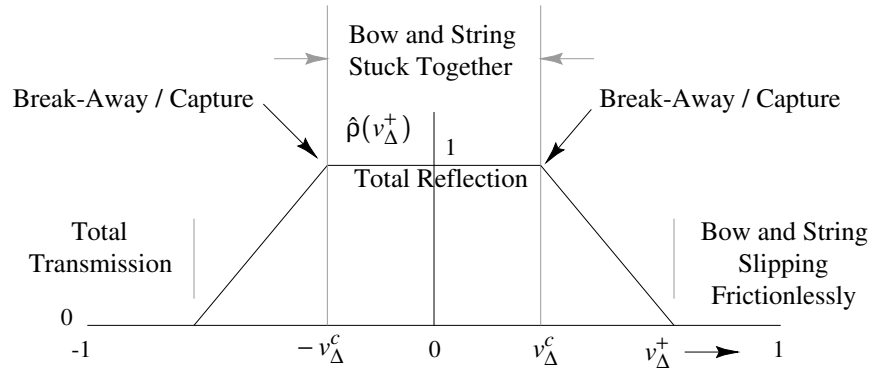
$v_s = v_{s,l}^+ + v_{s,l}^- = v_{s,r}^+ + v_{s,r}^- =$  transverse string velocity

$R_s =$  wave impedance of string

$R_b(v_\Delta) =$  friction coefficient for the bow against the string, i.e.,

$$F_b(v_\Delta) = R_b(v_\Delta) \cdot v_\Delta$$

## Simplified, Piecewise Linear Bow Table



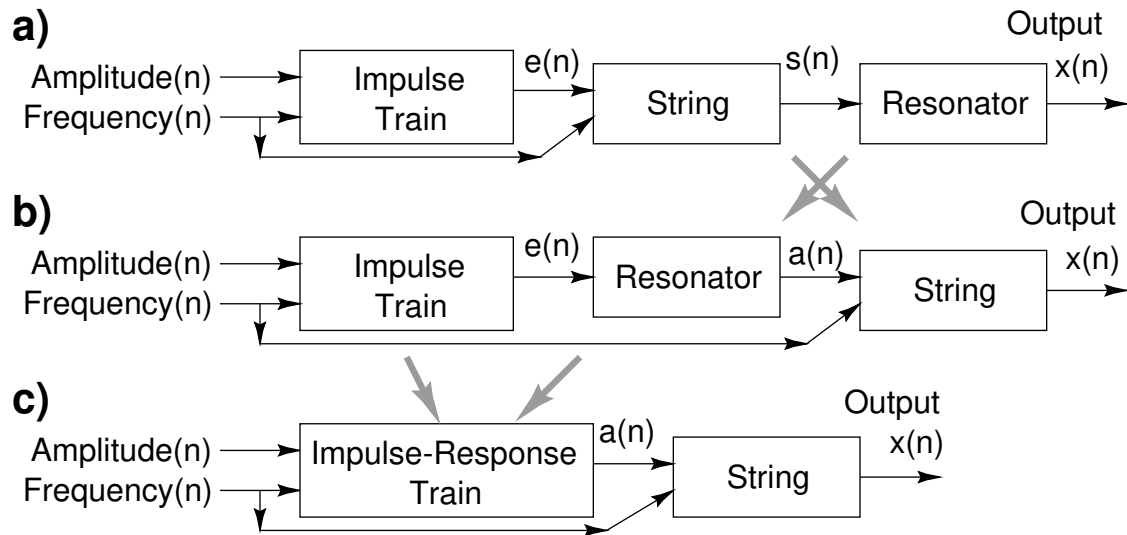
Simple, qualitatively chosen bow table for the digital waveguide violin.

- Flat center portion corresponds to a fixed reflection coefficient “seen” by a traveling wave encountering the bow stuck against the string
- Outer sections give a smaller reflection coefficient corresponding to the reduced bow-string interaction force while the string is slipping under the bow
- The notation  $v_{\Delta}^c$  at the corner point denotes the capture or break-away differential velocity
- Hysteresis is neglected



# Linear Commuted Violin Synthesis

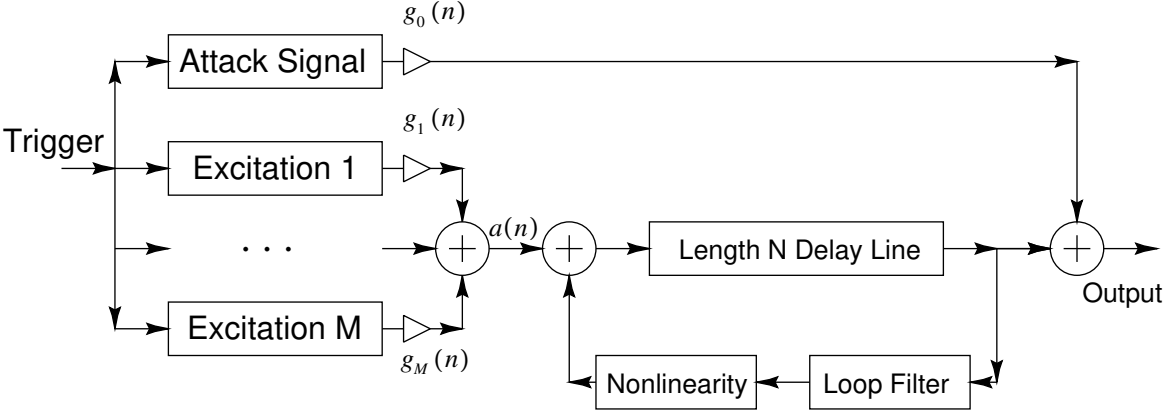
---



- Assumes *ideal Helmholtz motion*
- Sound examples:

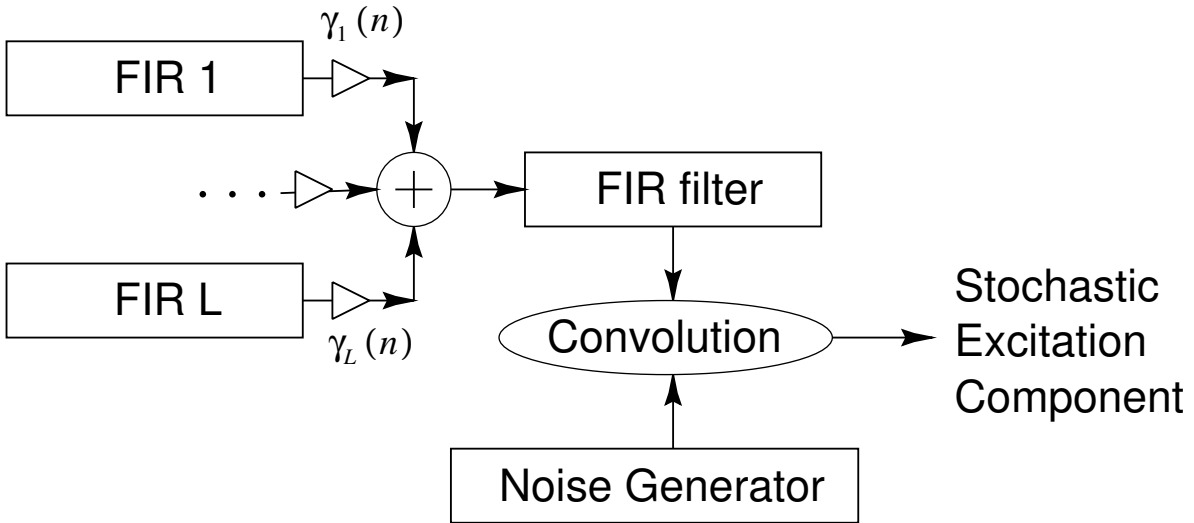
<http://ccrma.stanford.edu/~jos/wav/vln-lin-cs.wav>

# Multiple-Excitation Commuted Synthesis





# Filtered-Noise Excitation Synthesis



# Commuted Synthesis of the Linearized Violin

