

MUS420/EE367A Lecture 6D  
Plane-Wave Scattering at an Angle

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Outline

- Plane Wave at an Angle
- Plane-Wave Scattering at an Impedance Discontinuity
- Reflection and Refraction
- Evanescent Field due to Total Internal Reflection
- Plane-Wave Scattering at an Angle
- Imaginary wavenumbers

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- $\underline{u}$  = (unit) vector of *direction cosines*
- $k = \frac{2\pi}{\lambda}$  = (scalar) *wavenumber* along travel direction

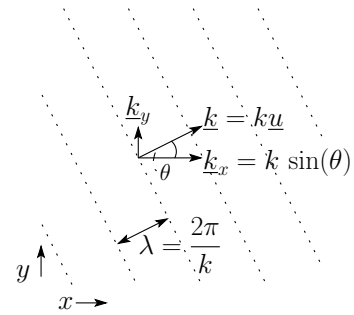
Thus, the vector wavenumber  $\underline{k} = k \underline{u}$  contains

- wavenumber in its magnitude  $k = \|\underline{k}\|$
- travel direction in its orientation  $\underline{u} = \underline{k}/k$

Note: wavenumber units are *radians per meter* (spatial radian frequency)

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Plane Wave at an Angle



wave crests of the sinusoidal traveling plane wave  
 $p(t, \underline{x}) = \cos(\omega t - \underline{k}^T \underline{x})$

Planar pressure-wave traveling in an arbitrary direction:

$$p(t, \underline{x}) = \cos(\omega t - \underline{k}^T \underline{x}), \quad \underline{x} \in \mathbf{R}^3$$

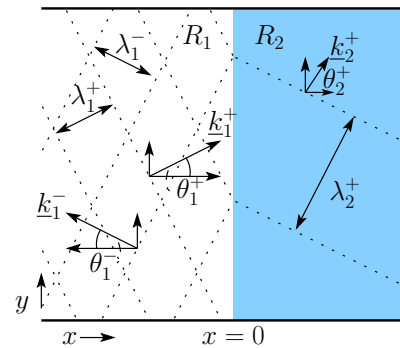
where  $\underline{k}$  = vector wavenumber:

$$\underline{k} = \begin{bmatrix} k_x \\ k_y \\ k_z \end{bmatrix} = k \begin{bmatrix} k_x/k \\ k_y/k \\ k_z/k \end{bmatrix} \triangleq k \begin{bmatrix} \cos \alpha \\ \cos \beta \\ \cos \gamma \end{bmatrix} \triangleq k \underline{u},$$

where

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Oblique Plane-Wave Scattering



By continuity, waves must agree on boundary plane:

$$\langle \underline{k}_1^+, \underline{r} \rangle = \langle \underline{k}_1^-, \underline{r} \rangle = \langle \underline{k}_2^+, \underline{r} \rangle$$

where  $\underline{r} = (0, y, z)$  denotes any vector in the boundary plane. Thus, at  $x = 0$  we have

$$k_{1y}^+ y + k_{1z}^+ z = k_{1y}^- y + k_{1z}^- z = k_{2y}^+ y + k_{2z}^+ z$$

If the incident wave is constant along  $z$ , then  $k_{1z}^+ = 0$ , requiring  $k_{1z}^- = k_{2z}^+ = 0$ , leaving

$$k_{1y}^+ y = k_{1y}^- y = k_{2y}^+ y$$

or

$$k_1 \sin(\theta_1^+) = k_1 \sin(\theta_1^-) = k_2 \sin(\theta_2^+)$$

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## Reflection and Refraction

Above we derived

$$k_1 \sin(\theta_1^+) = k_1 \sin(\theta_1^-) = k_2 \sin(\theta_2^+)$$

The first equality implies

$$\theta_1^+ = \theta_1^-$$

(Angle of incidence equals angle of reflection)

Let  $c_i$  denote the phase velocity in wave impedance  $R_i$ :

$$c_i = \frac{\omega}{k_i}, \quad i = 1, 2$$

In impedance  $R_2$ , we have in particular

$$\omega^2 = c_2^2 k_2^2 = c_2^2 [(k_{2x}^+)^2 + (k_{2y}^+)^2]$$

Solving for  $k_{2x}^+$  gives

$$k_{2x}^+ = \sqrt{\frac{\omega^2}{c_2^2} - (k_{2y}^+)^2} = \sqrt{\frac{\omega^2}{c_2^2} - k_2^2 \sin^2(\theta_2^+)}$$

Since  $k_1 \sin(\theta_1^+) = k_2 \sin(\theta_2^+)$  from above,

$$k_{2x}^+ = \sqrt{\frac{\omega^2}{c_2^2} - k_1^2 \sin^2(\theta_1^+)} = \sqrt{\frac{\omega^2}{c_2^2} - \frac{\omega^2}{c_1^2} \sin^2(\theta_1^+)}$$

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We have derived

$$k_{2x}^+ = \frac{\omega}{c_2} \sqrt{1 - \frac{c_2^2}{c_1^2} \sin^2(\theta_1^+)}$$

We earlier established  $k_{2y}^+ = k_{1y}^+$ .

- This describes the *refraction* of the plane wave as it passes through the impedance-change boundary.
- Refraction angle depends on ratio of phase velocities  $c_2/c_1$ .
- This ratio is often called the *index of refraction*:

$$n \triangleq \frac{c_2}{c_1}$$

and the relation  $k_1 \sin(\theta_1^+) = k_2 \sin(\theta_2^+)$  is called *Snell's Law* (of refraction).

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## Evanescent Field due to Total Internal Reflection

Note that if  $c_1 < c_2 |\sin(\theta_1^+)|$ , the horizontal component of the wavenumber in medium 2 becomes *imaginary*:

- Acoustic field in medium 2 is “*evanescent*”
- Wave in medium 1 undergoes “*total internal reflection*”
- No power travels from medium 1 to medium 2
- Evanescent field decays exponentially to the right
- “*Tunneling*” possible given medium 3 in which wave propagation resumes
- Reference: **Vibrations and Waves in Physics** by I.G. Main, Cambridge University Press, 1978.

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## What does it mean to have an imaginary wavenumber?

$$\begin{aligned} p(t, \underline{x}) &= \cos(\omega t - \underline{k}^T \underline{x}) = \operatorname{re}\{e^{j(\omega t - k_x x - k_y y)}\} \\ &= \operatorname{re}\{e^{j(\omega t - k_y y)} e^{-j k_x x}\}, \quad (\text{let } k_x \triangleq -j \kappa_x) \\ &\triangleq \operatorname{re}\{e^{j(\omega t - k_y y)} e^{-\kappa_x x}\} \\ &= e^{-\kappa_x x} \cos(\omega t - k_y y) \end{aligned}$$

- An imaginary wavenumber corresponds to an exponential decay
- Sign of  $\kappa_x$  is chosen to match boundary conditions at the plane
- Time dependence applies to *all* points to the right of the boundary (no “propagation”)

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