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A Real-Time System for Spatial Distribution of Sound

Marina Bosi

**CCRMA
DEPARTMENT OF MUSIC
Stanford University
Stanford, California 94305**

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Marina Bosi

Center for Computer Research in Music and Acoustics (CCRMA)

Department of Music, Stanford University

Stanford, California 94305

mab%ccrma-f4@sail.stanford.edu

Abstract: *One of the most significant parameters in the perception of music is related to the spatial characteristics of the sound. Composers have shown interest in exploitation of the spatial dimension of the sound since as early as the 16th century. We have precedents in such as works by Gabrieli, Berlioz, and, after the turn of the century, in Varese's Deserts, Stockhausen's Gesange der Juenglinge, Chowning's Turenas, etc...*

The goal of this work is the implementation of a tool which would allow the composer to project musical sound into an illusory acoustical space using a variable number of speakers. A two-dimensional model and the implementation for a single source and four channel transducers is presented. The user interface of the interactive software package, written in the C language, allows one to draw a sound path on the screen (which corresponds to the illusory space) and hear the displacement of the sound (complete with Doppler shift) in real time. Results and future work are discussed.

Sound Localization

In order for localization to occur, the listener needs to define the angular location of the sound source as well as the distance of the source relative to the listener. Our ability to localize sounds is influenced by many factors. On one side, the perception of sound direction based only on auditory information depends on the mechanism adopted by our brain to decode the sound stimuli (a striking example is the lateralization phenomenon). On the other side, the characteristics of the listening space (shape, size, absorption coefficient, etc..) and the characteristics of the sound source are directly related to the physics of the event. To give the illusion of spatial sound images, we need to bring the psychological and the physical aspects of the acoustical event into relation with the technology and the system used to reproduce the sound.

We will consider in this paper a two dimensional model, therefore the angular location is defined by azimuth (horizontal plane). From previous works on localization one

knows that the azimuthal localization cues, Interaural Time Difference (ITD) and Interaural Intensity Difference (IID), depend on the spectrum of the sound source (Fedderson et al. 1957): for wavelengths shorter than the interaural distance, interaural intensity differences are dominant cues, while for wavelengths longer than the interaural distance, interaural differences in arrival time of features of the sound waveform are dominant cues. It is possible to simulate the auditory system cues through headphones (Kendall and Rodgers 1981), however, since they are related to the precise location of the listener, trying to reproduce these cues in a concert situation, where the precise location of the listener is unpredictable, is unrealistic.

With regard to the distance cue, solving the wave equation in a free field, we know

Humidity	Frequency(Hz)				
	2000	3200	4000	5000	6400
40%	0.0032	0.0052	0.0072	0.0112	0.0172
50%	0.0028	0.0044	0.0060	0.0092	0.0144
60%	0.0028	0.0040	0.0056	0.0076	0.0120
70%	0.0028	0.0036	0.0052	0.0068	0.0100

Attenuation constant k of air (meter^{-1})

at 20⁰ C and normal atmospheric pressure

TABLE 1. Attenuation constant of air.

that the amplitude of the variation of the sound pressure is inversely proportional to the distance, r , between the source and the listener. Further, for large distances ($r \geq 100\text{m}$), we need to take into account the attenuation of the signal in the air which depends on the frequency of the source, as well as temperature, humidity and pressure (see Table1). This introduces an exponential decay of the form $e^{-k\tau}$. We have chosen to reproduce this data from Harris 1962 since erroneous data from Kuttruff 1973 has been widely distributed in the computer music field (e.g. Moorer 1979).

In enclosed environments intensity differences alone are only a relative cue to distance perception and it is the contribution of the reverberant energy that determines our ability to make distance discriminations (Sheeline 1982). The reverberant energy, which is

the the total energy minus the direct source energy, supplies the room information giving general cues as to size, shape and material construction . One might expect that the reflected sound represents a potential source of confusion in the localization of the sound source; however the reflected energy returned to the auditor within the following 50 ms is integrated with the direct energy into the impression of a single 'fuller', appropriately localized, sound. The accurate localization of a sound in spite of seemingly conflicting cues from reflections is sometimes called "law of the first wavefront" and sometimes the "precedence effect" (Haas 1951; Wallach et al. 1949).

When the sound energy is distributed between speaker pairs a sound image can be perceived at a different position than the physical sound sources position. If two loudspeakers with the same acoustic output are placed at equal but opposite angles from the listener on either side of the median plane, the listener perceives an apparent single image located in the midway position between the speakers. Moreover, if the intensity applied to the loudspeaker pair is not the same for each speaker, this can cause the illusion of an angular displacement in the sound image (Gardner 1973).

The Model

Our goal is to attain control over the spatial characteristics of musical sound through a system which allows real-time interaction and which would be portable and flexible enough to be used in many different situations (Moore 1983). We need a general way to obtain reasonable-sounding values for the processing algorithm. The model described here is based on the following elements:

- The position of the sound source within the illusory space relative to the listener.
- The number of loudspeakers and their locations as well as the characteristics of the listening space such as shape, size, absorption, etc.

The configuration of the variable number of loudspeakers can be arbitrarily chosen in order to reproduce any possible "real" configuration (for example an open air concert situation where the loudspeakers are placed in semicircle in front of the public, or the classical quad configuration, or any other precise location one needs).

Let's consider the the system configuration of figure 1. We want to simulate a source located, at a certain instant, at the point S. We choose to represent the point S in polar coordinates (r, ϑ) where the origin is at the listener L and ϑ is the angle relative to the speaker 1. The listener is surrounded by N loudspeakers numbered counter-clockwise. The location of each speaker i can be represented by the doublet (R_i, Θ_i) .

We know that, in a free field for small distances, the intensity I relative to a sound source which irradiates omnidirectionally is

$$I = \frac{I_0}{r^2}$$

where I_0 is the intensity at the distance $r_0 \equiv 1$. For $r \leq r_0$ our system will deliver the maximum sound intensity level, corresponding to the intensity I_0 . For large distances we need to consider the additional factor in the dependence of the intensity on the distance r .

$$I = \frac{I_0}{r^2} e^{-kr}$$

where the constant k depends on the absorption of sound in the air. At fixed humidity, pressure and temperature values, the absorption coefficient k increases with increasing frequency of the signal. The absorption in the air, whose effect is negligible for small distances ($< 100\text{m}$ from table 1), acts like a low pass filter, highly attenuating the high frequencies for large distances.

As we noted before, sound parameters dependent on interaural time differences, hence dependent on the precise position and orientation of the listener head, are uncontrollable when music is produced through a multiple channel loudspeakers system. Each adjacent speaker pair, i and $i+1$, defines a portion of the plane between the directions Θ_i and Θ_{i+1} . When the source is located in that portion of the plane, we can simulate the angular location by distributing the source intensity between the speaker pair and applying a zero intensity to all the other speakers. The intensity coming from each speaker direction is determined by the ratio of the angular position of the source relative to the speakers, and the angular separation of the two speakers (Chowning 1971). The fraction of the intensity supplied by each speaker pair $i, i+1$ will be

$$p_i = \frac{\Theta_{i+1} - \vartheta}{\Theta_{i+1} - \Theta_i} = 1 - p_{i+1} \quad \Theta_i < \vartheta < \Theta_{i+1} \quad i = 1, 2, \dots, N-1$$

$$p_N = \frac{2\pi - \vartheta}{2\pi - \Theta_N} = 1 - p_1 \quad \Theta_N < \vartheta < 2\pi \quad i = N$$

where (R_i, Θ_i) defines the position of the loudspeaker i and $\Theta_1 \equiv 0$.

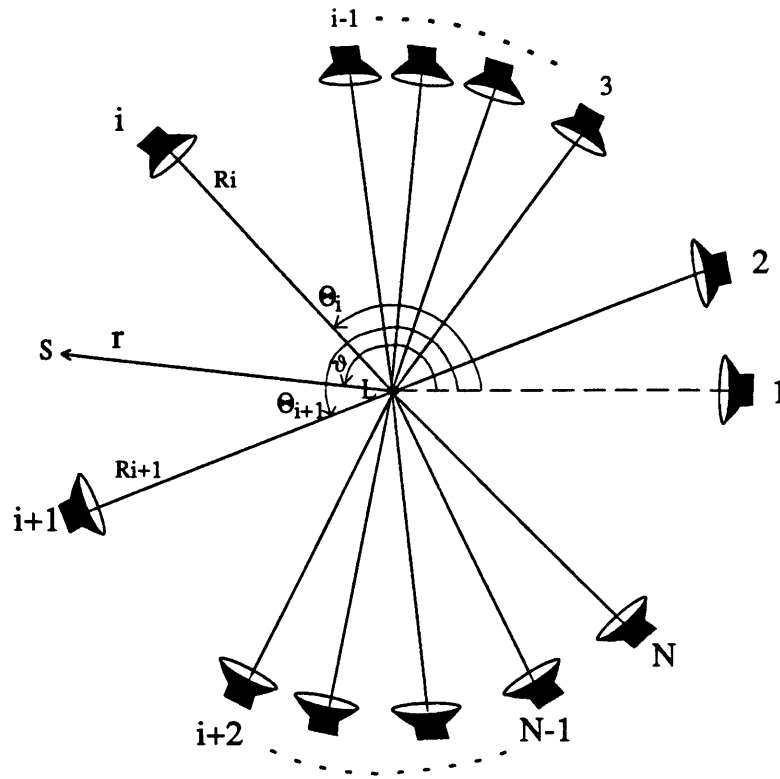


FIG. 1. Loudspeakers configuration

To take into account that the speakers in general will be at different distances R_i from the listener, we also need to multiply the source intensity by the factor

$$C_i = \left(\frac{R_i}{R_{\max}}\right)^2$$

Notice if a speaker is farther away than another we make that speaker louder to correct for it, and vice-versa.

A variation in the position, r , of the source will affect its frequency through the Doppler effect. We need to consider the corrected frequency

$$f' = f \frac{c}{c + \hat{\mathbf{r}} \cdot \vec{v}_s}$$

where f is the frequency of the motionless source, c is the speed of the sound in the air, $\hat{\mathbf{r}}$ is the unit vector in the radial direction and $\vec{v}_s = \frac{d\vec{r}}{dt}$ is the velocity of the sound source. We choose the origin of our reference system where the 'ideal' listener is located, therefore the

component of the source velocity along the unit vector $\hat{r} \equiv \frac{\vec{r}}{r}$ will be negative when the source moves toward the listener (i.e. the frequency will increase) and positive when the source moves away (i.e. the frequency will decrease).

To summarize, given the position of the speakers, (R_i, Θ_i) , and, at each instant, the position of the source, (r, ϑ) and its velocity, \vec{v}_s , we can determine at each instant the corrected frequency, f , and the relative intensity values delivered by the active channels i and $i+1$:

$$I_i = I C_i p_i = \frac{I_0}{r^2} e^{-kr} \left(\frac{R_i}{R_{\max}} \right)^2 \frac{\Theta_{i+1} - \vartheta}{\Theta_{i+1} - \Theta_i}$$

$$I_{i+1} = I C_{i+1} p_{i+1} = \frac{I_0}{r^2} e^{-kr} \left(\frac{R_{i+1}}{R_{\max}} \right)^2 \frac{\vartheta - \Theta_i}{\Theta_{i+1} - \Theta_i}$$

We may easily extend this algorithm to the case of more than one sound source. Let's say we have $j = 1, 2, \dots, M-1, M$ sources represented by the points (r_j, ϑ_j) . We would apply the results above to each source and superimpose at each instant the intensities $I_{j,i}$ and $I_{j,i+1}$ for each channel pair, i and $i+1$, activated by every source j .

The reverberation plays an essential role in the perception of distance. In our model we consider the reverberant energy to be spatially uniform, therefore the ratio of reverberant to direct sound energy reaching the listener will increase as the distance r between the sound source and the listener increases.

Implementation

In figure 2 is shown the system implementation for a single source ($j = 1$) and four channel transducers ($i = 1, 2, \dots, 4$). A Macintosh II sends MIDI control signals to a Yamaha DMP-7 digital mixer which provides modules for the spatial processing algorithm. MIDI control signals are sent as well to a synthesizer source (Yamaha DX7 at the present). Once the position of the speakers is selected, as the user draws a sound trajectory on the screen with the mouse, the information about the position (r, ϑ) of the sound source with respect to the listener and the velocity \vec{v}_s are detected. The intensity of the source, $I(r)$, is

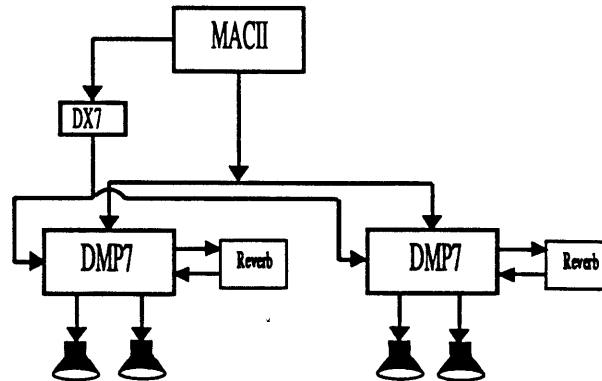


FIG. 2. System schematic

computed and the active channel pair is determined based on the angular location of the source. We now can determine the normalized intensity supplied by each speaker in the active channel pair :

$$I_i = [I C_i p_i] / I_0$$

$$I_{i+1} = [I C_{i+1} p_{i+1}] / I_0 .$$

where I_0 is the maximum intensity (i.e. the intensity when the distance between the source and the listener is $r = r_0 = 1$), C_i is the factor that takes into account differences in distance between speakers and p_i is the panning factor .

The value I' for each speaker is then mapped to a midi number in the range 0-127 according to the DMP7 fader control function. We have measured this function and found that the intensity is quadratically related to the midi number for midi numbers in the range [0,50]. Beyond this range there is no such simple relation and so we resort to a polynomial fit. We find cubic order (figure 3) is enough.

The velocity of the sound source determines the Doppler correction factor g to the value of its frequency, where $g = c / (c + \hat{\mathbf{r}} \cdot \vec{\mathbf{v}}_s)$. If we limit the speed of the sound source $\hat{\mathbf{r}} \cdot \vec{\mathbf{v}}_s$ in the range $[0, c/2]$, then the frequency of the sound source varies in the range $[f - f/2, f + 2f]$, (i.e. plus or minus an octave). This is mapped to a midi number affecting the pitch bend function of the synthesizer (Loy 1985).

The reverb module used (the DMP-7 rev room effect at the present) allows one to modify the reverberation time and to set an initial delay which depends both on the source position and the listening space shape. A control over the early reflection patterns which depends on the size and shape of the room and the source position is also implemented.

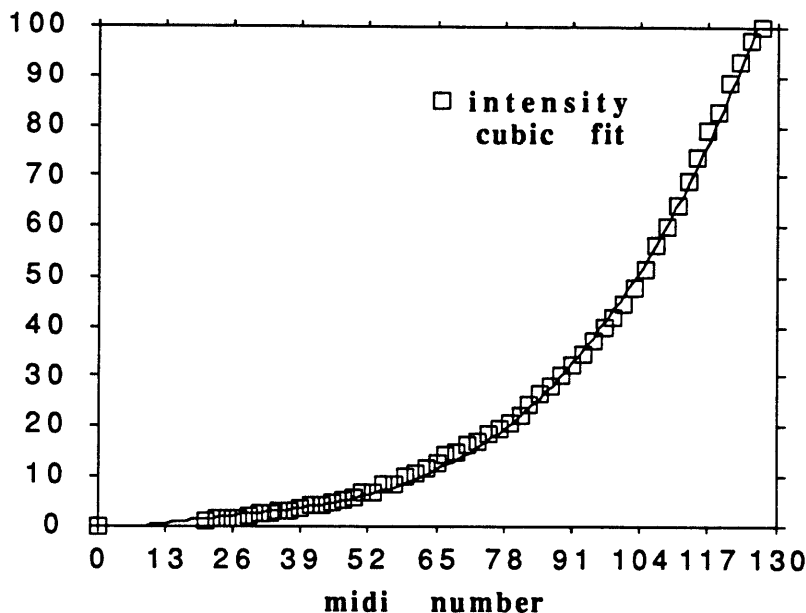


FIG. 3. Measured correspondence between midi numbers and intensity (normalized values).

The user interface is shown in figure 4. As one draws a sound path on the screen (which corresponds to the illusory space) one can hear the displacement of the sound in the acoustical environment specified by the user. The control panel allows the user to select the amount of reverberation, the room size, as well as the display scale. The source position (x and y coordinates in meters) within the illusory space is indicated at each instant at the bottom of the control panel. It is possible to record sound trajectories and to create and edit files of recorded trajectories. The loudspeaker configuration within the space can be varied by pressing the option key while dragging the loudspeakers (indicated by black squares on the screen) with the mouse.

Results/Future

The goal of this work was to provide the composer with a tool to project the musical sound into an illusory acoustical space using a variable number of speakers. The model implementation, realized in a real-time implementation for the Macintosh II and MIDI devices, allows one, through an interactive software package written in the C language, to draw a sound trajectory on the screen and hear the movement of the sound (complete with

Doppler shift) in real time. This system has been tested in different acoustical spaces using a quadraphonic reproduction system and listeners found that it produces convincingly localized sound images.

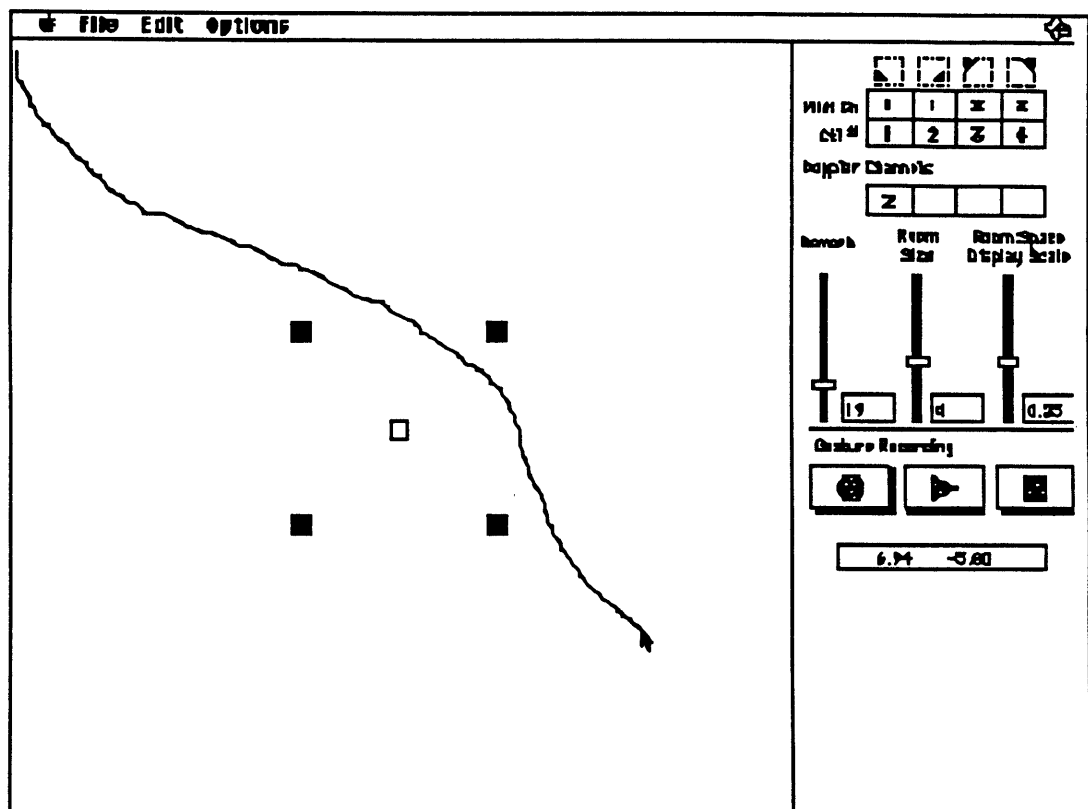


FIG. 4. User interface.

Although the initial implementation allows for simple room geometries, one of our goals for future work is to take more accurate account of the features of the acoustical space. For instance we would like to include multiple ray paths for each source. We also would like to replace the use of the digital mixer with a spatial processing software for a specialized DSP chip like the Motorola 56000. The model presented could readily be extended to three dimension.

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